

# Performance Dependence on Doping of THz Quantum-Cascade Lasers

Aaron Maxwell Andrews, Alexander Benz, Gernot Fasching, Karl Unterrainer, Tomas Roch, Werner Schrenk, and Gottfried Strasser

**Abstract**—We present the effects of dopant concentration on the performance of 2.75 THz quantum-cascade lasers. The sheet density was varied from  $5.4e9$  to  $1.9e10$   $\text{cm}^{-2}$  by growing 4 identical LO-phonon depletion scheme structures and additionally an undoped structure. The threshold current density  $J_{\text{th}}$  and maximum current density  $J_{\text{max}}$  scaled linearly with doping concentration while the applied field for the onset of lasing and cessation of lasing remained independent. The maximum operating temperature for all structures was  $\sim 140\text{K}$ . These results suggest that within this doping range free carrier absorption is not limiting performance for THz quantum-cascade lasers. Additionally, an undoped structure is expected to have a  $J_{\text{th}}$  of zero and a  $J_{\text{max}}$  of  $70$   $\text{A}/\text{cm}^2$ .

**Index Terms**— Doping in Lasers, Quantum-Cascade Lasers, Semiconductor lasers, Terahertz lasers.

## I. INTRODUCTION

With the realization of the first THz quantum-cascade laser (QCL) [1] there has been great progress in available emission frequencies down to 1.6 THz [2], reduction of threshold current  $J_{\text{th}}$  to  $1$   $\text{A}/\text{cm}^2$  [3], and maximum operating temperature up to  $164\text{K}$  [4]. This has been accomplished through the improvement of QCL active region and waveguide designs: chirped superlattice [1] and bound-to-continuum [5] (low threshold designs), LO-phonon depopulation scheme [6] (high-temperature design), surface plasmon and double metal waveguides.

A previous study into the doping characteristics of the LO-phonon THz QCL structures [7] used a delta doping gradient, created by stopping sample rotation, producing within the same sample a doping variation of  $3\text{-}5e10$   $\text{cm}^{-2}$  or  $\pm 20\%$ . In

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this paper, we investigate four uniformly doped samples over a wide range with approximately  $+50\%$  differences between doping concentrations, including an additional undoped THz QCL.

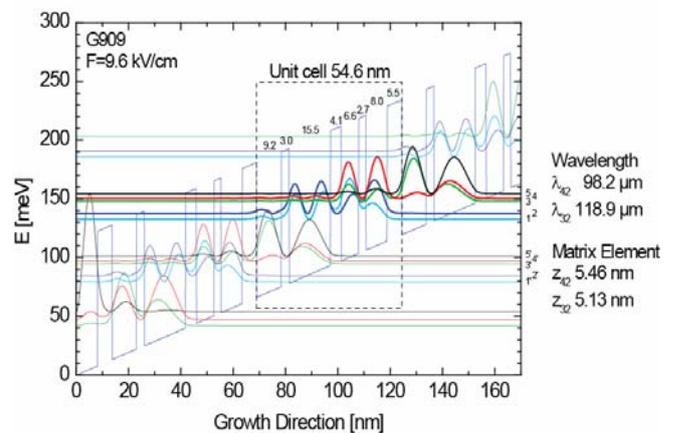


Fig. 1. Calculated band structure at a field of  $9.8$   $\text{kV}/\text{cm}$ . The box denotes one unit cell of the  $15$   $\mu\text{m}$  active region. The barriers are  $\text{Al}_{0.15}\text{Ga}_{0.85}\text{As}$  and the wells are  $\text{GaAs}$ . Lasing is between levels 4 and 2 and 3 and 2, depending on the applied field. The lower laser state is depopulated into the ground states  $5'$ ,  $4'$ , and  $3'$  of the next cascade by an LO phonon.

## II. EXPERIMENT

Five LO-phonon depopulation 2.75 THz QCL structures were grown by solid-source molecular-beam epitaxy (MBE) as in Ref. [8], shown in Fig. 1, with 271 cascade cells for a  $15$   $\mu\text{m}$  thick active region to minimize waveguide losses. This design uses  $\text{Al}_{0.15}\text{Ga}_{0.85}\text{As}$  barriers with  $\text{GaAs}$  wells to define the upper and lower laser states as well as an LO-phonon resonance with the lower laser level to rapidly deplete the electrons and prevent thermal back filling. The doping in the widest well in the five samples was  $5.4e9$ ,  $8.2e9$ ,  $1.2e10$ ,  $1.9e10$   $\text{cm}^{-2}$ , and undoped. The accuracy of the structure was determined by high-resolution x-ray diffraction. The samples were processed with an Au-Au double metal waveguide for high mode confinement and the laser sidewalls and facets were etched, instead of cleaved, by inductively coupled reactive ion etching (ICP-RIE). The dimensions for each laser bars are  $120\mu\text{m} \times 1\text{mm}$ .

Device measurement was performed in pulsed mode

operation with a pulse length of 100 ns, chosen to prevent additional device heating, and a repetition rate of 1 kHz for the integral measurements and 5 kHz for spectral

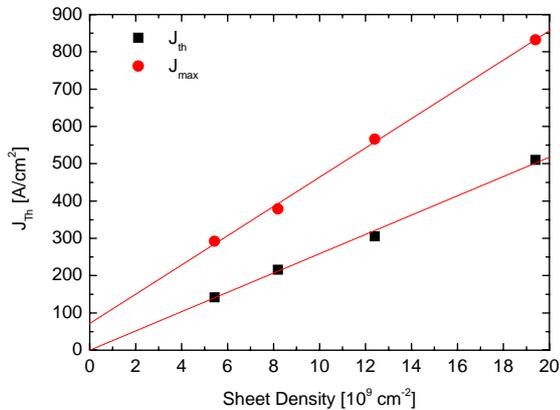


Fig. 2. Plot of threshold current density  $J_{th}$  and maximum current density  $J_{max}$  vs doping sheet density for the four doped structures. Both  $J_{th}$  and  $J_{max}$  show a linear dependence on doping density.

measurements.

### III. RESULTS AND CONCLUSIONS

The undoped sample shows similar electrical characteristics as the four doped samples. The band structure aligns at 9 kV/cm and then misaligns at a current density of 60 A/cm<sup>2</sup>, just below the predicted 70 A/cm<sup>2</sup> from the linear fit to the doped data points. This may indicate a lower doping or impurity limit of the MBE or in a naive approach this corresponds to an overall background doping of 1-2e14 cm<sup>-3</sup> for the 22 hour growth.

Lasing, except for the undoped sample, in the structures starts at an applied field of 8 kV/cm, where the alignment of the cascades begins and the I-V has a reduced slope. The threshold current density  $J_{th}$  increases linearly with carrier concentration. The maximum current density  $J_{max}$  also increases linearly with carrier concentration. Thus the carrier concentration also increases the operating range of the laser. While the  $J_{th}$  converges on zero A/cm<sup>2</sup>, the  $J_{max}$  converges on 70 A/cm<sup>2</sup>. This is in contrast with Ref. [7], where non-linear behavior was observed for the similar LO-phonon design with higher dopant concentration. The non-linearity was attributed to free-carrier absorption and a strong dependence of  $T_{max}$  on carrier density. The  $T_{max}$  was ~145K for all samples in our study.

The linearity of the  $J_{th}$  and  $J_{max}$  with respect to doping and the independent nature of  $T_{max}$  leads to the conclusion that free carrier absorption and/or impurity scattering is not yet limiting performance in these THz QCLs. The doping can be changed to directly produce lower  $J_{th}$ , although lower doping results in lower output power.

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